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NUMERICAL SOLUTION OF LINEAR
PARABOLIC PARTIAL DIFFERENTIAL
EQUATION WITH INFINITE RANGE
FOR SPACE VARIABLE

By

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Abstract

Linear parabolic partial differential equation is encountered in many areas of science. Some of the well-known examples are heat conduction in solids, dispersion of tracers in geophysical reservoirs, movement of groundwater and moisture in soils, dispersal of pollutants in water and air etc.

We have used the well-known Crank-Nicholson Scheme in solving a general linear parabolic partial differential equation with infinite range for space variable. A new feature of the algorithm developed by us is the transformation of the space variable to make the spatial grid equally spaced between zero and infinity. This procedure saves efforts as well as computer time.

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This report deals with the development of finite difference method of solving a linear parabolic partial differential equation of the form:

$$\frac{\partial u}{\partial t} = f_1(t, y) \frac{\partial^2 u}{\partial y^2} + f_2(t, y) \frac{\partial u}{\partial y} + f_3(t, y) u + f_4(t, y) \dots \quad (1)$$

with the initial and boundary conditions as:

$$\begin{aligned} u(0, t) &= g_1(t) &&) \\ u(a_0, t) &= g_2(t) &&) \\ u(y, 0) &= g_3(y) &&) \end{aligned} \dots \quad (2)$$

where f_1, f_2, f_3, f_4 are specified functions of t and y ; g_1, g_2 are given functions of t only and g_3 is a given function of y .

For finite difference methods, it is known that step size in the spatial direction has to be smaller when the gradients are higher. Thus step size should be smaller near $y = 0$ and has to be made progressively larger as $y \rightarrow \infty$. Such a choice of grid making, however, creates practical difficulties. To obviate the above difficulty, we introduce the transformation

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$$z = \exp(-y^\beta) \dots (3)$$

which brings the range of spatial variable z between 1 and 0 corresponding to the range of y variable between 0 and ∞ .

This transformation has the additional advantage of uniform spacing of the interval (0, 1) in the z -variable.

The grid points in the z variable are 1, $(n-1)h$, $(n-2)h$, ----- $2h$, h , 0; where $h = 1/n$ (zero corresponds to ∞ of the original spatial variable y). The choice of β can be made in such a way that

$$y = (-\ln z)^{1/\beta}$$

is of the order of infinity for the problem in question.

From (3)

$$z' = \frac{dz}{dy} = z (-\beta y^{\beta-1}) \dots (4)$$

$$z'' = \frac{d^2z}{dy^2} = z y^{\beta-1} \left[\beta y^{\beta-1} - (\beta-1) \frac{1}{y} \right] \dots (5)$$

Also

$$\frac{\partial u}{\partial y} = \frac{\partial u}{\partial z} \frac{dz}{dy} = \frac{\partial u}{\partial z} z' \dots (6)$$

$$\frac{\partial^2 u}{\partial y^2} = \frac{\partial u}{\partial z} \frac{d^2z}{dy^2} + \frac{\partial^2 u}{\partial z^2} \left(\frac{dz}{dy} \right)^2 = \frac{\partial u}{\partial z} z'' + \frac{\partial^2 u}{\partial z^2} z'^2 \dots (7)$$

using (6) and (7) the given parabolic differential equation,

(1), takes the form

$$\frac{\partial u}{\partial t} = f_1(t, z) \left[z'^2 \frac{\partial^2 u}{\partial z^2} + z'' \frac{\partial u}{\partial z} \right] + f_2(t, z) z' \frac{\partial u}{\partial z} + f_3 u + f_4$$

$$\text{or, } \frac{\partial u}{\partial t} = f_1 z'^2 \frac{\partial^2 u}{\partial z^2} + (f_1 z'' + f_2 z') \frac{\partial u}{\partial z} + f_3 u + f_4 \dots (8)$$

with the transformed initial and boundary conditions corresponding to (2) as:

$$\begin{aligned} u(i, t) &= g_1(t) \\ u(0, t) &= g_2(t) \\ u(z, 0) &= g_3(z) \end{aligned} \quad \dots (9)$$

Applying the Crank-Nicholson scheme, the finite difference at the i^{th} space interval and for the time interval between j and $j + 1$ is given as:

$$\frac{u_i^{j+1} - u_i^j}{\Delta t} = \left[(E_i)^{j+1} + (E_i)^j \right] / 2 \quad \dots (10)$$

where

$$t_j = j \cdot \Delta t \quad \dots (11)$$

$$E_i = f_{1i} z_i^2 \frac{u_{i+1} - 2u_i + u_{i-1}}{h^2} + (f_{1i} z_i + f_{2i} z_i)$$

$$\frac{u_{i+1} - u_{i-1}}{2h} + f_{3i} u_i + f_{4i} \quad \dots (12)$$

In (12) $f_{1i} = f_1(t, z_i)$ and likewise f_{2i} , f_{3i} and f_{4i} .

Writing out explicitly the terms using (11) and (12), the finite-difference equation (10) takes the form

$$A_{i-1} u_i^{j+1} + B_{i-1} u_{i+1}^{j+1} + F_{i-1} u_{i-1}^{j+1} = d_{i-1};$$

$$i = 2, 3, \dots, n \quad \dots (13)$$

where

$$\begin{aligned}
 A_{i-1} &= \frac{2}{\Delta t} + \frac{2 f_{1i}^{j+1}}{h^2} z_i'^2 f_{3i}^{j+1} \\
 B_{i-1} &= -f_{1i}^{j+1} z_i'^2/h^2 - (f_{1i}^{j+1} z_i'' + f_{2i}^{j+1} z_i')/2h \\
 F_{i-1} &= -f_{1i}^{j+1} \frac{z_i'^2}{h^2} + (f_{1i}^{j+1} z_i'' + f_{2i}^{j+1} z_i')/2h \\
 d_{i-1} &= (E_i)^j + f_{4i}^{j+1} + \frac{2}{\Delta t} u_i^j
 \end{aligned}
 \quad \dots (14)$$

Equation (13) represents (n-1) simultaneous linear equations in (n-1) unknowns at (j+1)th level of the variable t, namely

$$\begin{bmatrix}
 A_1 & B_1 & 0 & 0 & \dots & 0 & 0 & 0 \\
 F_1 & A_2 & B_2 & 0 & \dots & 0 & 0 & 0 \\
 0 & F_2 & A_3 & B_3 & \dots & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & \dots & F_{n-3} & A_{n-2} & B_{n-2} \\
 0 & 0 & 0 & 0 & \dots & 0 & F_{n-2} & A_{n-1}
 \end{bmatrix}
 \begin{bmatrix}
 u_2 \\
 u_3 \\
 \vdots \\
 \vdots \\
 u_n
 \end{bmatrix}
 =
 \begin{bmatrix}
 d_1 \\
 d_2 \\
 \vdots \\
 \vdots \\
 d_{n-1}
 \end{bmatrix}
 \quad \dots (15)$$

; i = 1, 2 ----- n

The boundary conditions will alter the first and last component of the right hand side of (15) as:

$$\begin{aligned}
 d_1 &= d_1 - F_1 u_1 \\
 d_{n-1} &= d_{n-1} - B_{n-1} u_{n+1}
 \end{aligned}
 \quad \dots (16)$$

where u_1 and u_{n+1} are known from the given boundary conditions.

For $t = 0$, the profile of u is given (i.e., initial condition), the profile at $t = \Delta t$ can be obtained by solving these simultaneous equations. We have solved this tridiagonal system of linear

equations by the Gaussian Elimination Method (Subroutine "TRID", Appendix 1). This subroutine returns the solution in variables F_1 to F_{n-1} . The values of u at the required points of the space variable y has been done using the nearest four points in the solution, (two above and two below) by fitting a third degree polynomial.

It may be noted that in some physical problems the initial condition

$$u(y, 0) = g(y)$$

may be the form

$$u(y, \infty) = g(y)$$

In such cases the transformation $T = 1/t$ will bring the problem to the above formulation.

AN ALGORITHM FOR THE FINITE DIFFERENCE SCHEME

In order to illustrate the above method we give an algorithm (Appendix 1) for solving the second order partial differential equation:

$$\frac{\partial c(x, t)}{\partial t} = D \frac{\partial^2 c}{\partial x^2} - u \frac{\partial c}{\partial x} \quad \dots (17)$$

For the initial and boundary conditions:

$$\begin{array}{l} c(0, t) = C_0 \sin \omega t \text{ -----(a) } \\ c(\infty, t) = 0 \text{ -----(b) } \\ c(x, 0) = 0 \text{ -----(c) } \end{array} \quad \dots (18)$$

Equations (17) and (18) describe the behaviour of a tracer in a dispersive groundwater system. 'C' represents the tracer concentration (dimension, ML^{-3}) at a time 't' at a distance 'x' from the recharge boundary ($x=0$) of the aquifer. Thus it is assumed that the tracer is injected continuously into the aquifer at the recharge boundary, the concentration varying with time sinusoidally as given by equation 18(a). The tracer flows with the water with a velocity 'u' and during the flow the tracer gets dispersed into the aquifer, the dispersion coefficient being D (dimension, $L^2 T^{-1}$). It is customary to use non-dimensional distance parameter xu/D (say 'y' for brevity) and the dispersion time constant, D/u^2 (dimension, T^{-1}). With these notations (17) becomes

$$\frac{\partial c(y,t)}{\partial t} = \frac{1}{(D/u^2)} \frac{\partial^2 c(y,t)}{\partial y^2} - \frac{1}{(D/u^2)} \frac{\partial c(y,t)}{\partial y} \dots (19)$$

with the transformed initial and boundary conditions as:

$$C(y,0) = 0; C(0,t) = C_0 \sin \omega t; C(\infty,0) = 0 \dots (20)$$

On comparing (19) and (20) with the general linear parabolic partial differential equation (1) and the initial and boundary conditions (2) respectively, one obtains

$$\begin{aligned} f_1(t,y) &= 1/(D/u^2) &&) \\ f_2(t,y) &= -1/(D/u^2) &&) \\ f_3(t,y) &= 0 &&) \\ f_4(t,y) &= 0 &&) \end{aligned} \dots (21)$$

$$\begin{aligned}
 g_1(t) &= C_0 \sin(\omega t) \\
 g_2(t) &= 0 \\
 g_3(y) &= 0
 \end{aligned}
 \quad \dots (22)$$

Equation (19) has been solved by the Crank-Nicholson Method for the given initial and boundary conditions (18) by appropriately defining (using equations 21 and 22) $f_1, f_2, f_3, f_4, g_1, g_2$ and g_3 in FUNCTION Subprograms F1, F2, F3, F4, G1, G2, and G3 given in the algorithm (Appendix 1). In order to get the normalized concentrations (i.e. $C(y, t)/C_0$) we put $C_0 = 1$ wherever it appears in (21) and (22).

In Fig. 1 the continuous curves A and B give $C(y, t)/C_0$ vs. t at $xu/D = 5$ and 10 respectively, obtained from the method given above. The value of D/u^2 taken for both the curves is 100 years and that of $\omega = 2\pi/T = 2\pi/1000 \text{ yr}^{-1}$ (i.e. period of the sine wave as 1000 years, see inset in Fig. 1).

We have also solved equation (17) for the initial and boundary conditions given in (18) analytically; Appendix 2 gives the analytical solution. The values obtained from the analytical solution for the same values of parameters D/u^2 and xu/D are shown by 'dots' in Fig. 1. It is seen that the numerical solution agrees closely with the analytical solution.

Table 1 gives a comparison of the values of tracer concentration obtained from the numerical and analytical methods at $xu/D = 5$ and for various values of time (normalised with the period of the sine wave). The error is the difference of the solutions

expressed as a percentage of the analytical value.

t/T	Finite Difference Solution	Analytical Solution	Difference	Percentage Error
0.25	0.067171	0.067277	-0.000106	-1.57
0.50	0.459601	0.459920	-0.000319	-0.07
0.75	0.341715	0.341093	+0.000622	0.18
1.0	-0.270253	-0.270703	+0.000450	-0.17
1.5	0.305339	0.305933	-0.000594	-0.19
2.0	-0.298452	-0.298945	+0.000493	-0.16

The Crank-Nicholson Implicit Method has advantage over the "Explicit" methods in that the former is unconditionally stable and convergent for all finite values of step size of both space and time variables. In addition for a linear problem, the resulting tridiagonal system of equations can be conveniently solved saving a lot of computer time.

Acknowledgements

We thank Shri Avinash Khare for help in obtaining the analytical solution of dispersion convection equation given in Appendix 2.

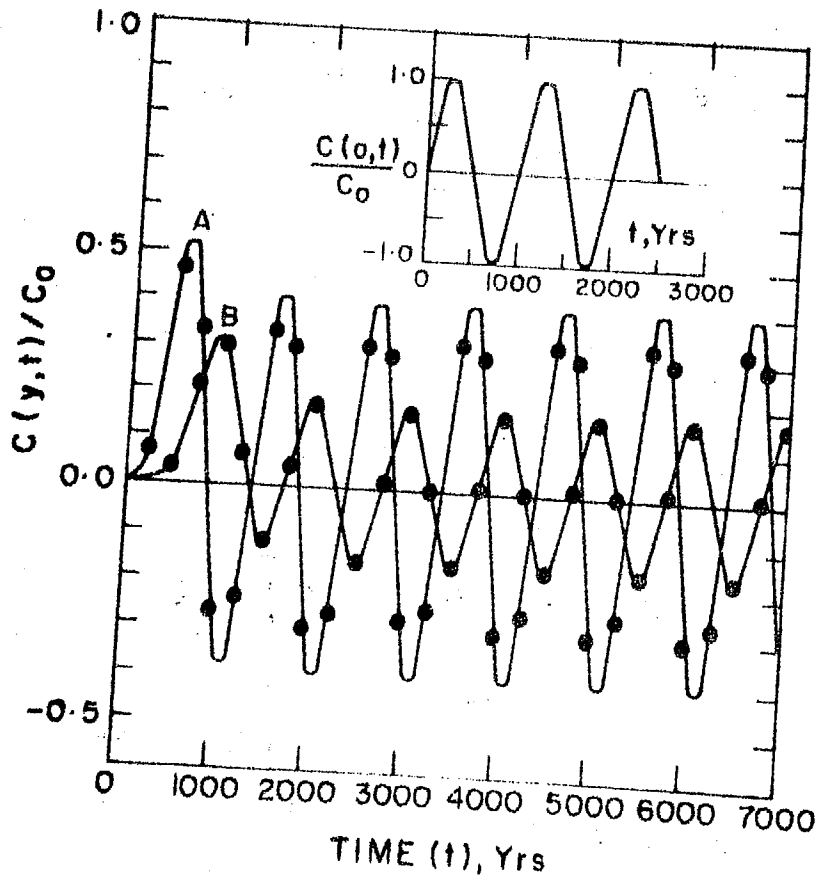


Fig. 1: Normalized radiotracer concentration, $C(y,t)/C_0$, in an aquifer having sinusoidal tracer concentration input at the recharge boundary ($y=0$) as shown in the inset. The value of D/u^2 is 100 years; the value of xu/D is 5 for curve A and 10 for curve B. Dots represent values obtained from the analytical solution.

APPENDIX 1

PROGRAM SOLVES A GENERAL LINEAR PARTIAL DIFFERENTIAL
 PARABOLIC EQUATION

$$DU/DT = F1(T,Y)*DU/DY^2 + F2(T,Y)*DU/DY + F3(T,Y)*U + F4(Y)$$
 WHERE DU/DT IS PARTIAL DERIVATIVE OF U WITH RESPECT TO T
 DU/DY IS PARTIAL DERIVATIVE OF U WITH RESPECT TO Y AND
 D²U/DY² IS PARTIAL SECOND DERIVATIVE OF U WITH RESPECT TO Y
 WITH THE INITIAL AND BOUNDARY CONDITIONS

U(0,T) = G1(T)
 U(INFINITY,T) = G2(T), AND
 U(Y,0) = G3(Y)

F1,F2,F3,F4 ARE SPECIFIED FUNCTIONS OF T AND Y; G1 AND G2 ARE
 GIVEN FUNCTIONS OF T AND G3 IS FUNCTION OF Y
 FUNCTIONS F1,F2,F3,F4,G1,G2 AND G3 GIVEN HERE ARE FOR THE CASE
 DISCUSSED IN PAGE 7 OF THE TEXT (SEE FIG. 1)
 THE USER HAS TO DEFINE THESE FUNCTIONS AS PER NEED

IMPLICIT REAL*8(A-H,C-Z)

REAL*4DENS

DIMENSION Y(1000),TIMES(1000),Z(1000),TZ(1000),CZ(1000),C(1000),
 CA(1000),B(1000),D(1000),E(1000),UU(1000),YC(100),DENS(100,200)
 COMMON AM,AM,LM,NI,N,K1,K2,K3,K4

IF=0

LM=C

N=100

NI=N+1

AM=N-1

IF K1=1 : 1=1/T VALUES WILL BE TAKEN FOR PRINTING AS WELL AS
 PLOTTING

IF K1=0 : NO INVERSE TRANSFORMATION BUT LINEAR TIME VALUES WILL
 BE TAKEN FOR PRINTING AND PLOTTING

IF K2=1 : LOG10 VALUES OF X AND Y WILL BE TAKEN FOR PRINTING
 AND PLOTTING

IF K2=0 : X AND Y VALUES AS SUCH TAKEN (LINEAR PLOT)

IF K3=1 : PLOTS REQUIRED ; K3=0 : NO PLOTS REQUIRED

IF K4 = 1 : X AND Y VALUES AT INTERPOLATED 'YC' VALUES ARE PRINTED

READ 200,K1,K2,K3,K4

2((FORMAT(13)

TIME=FINAL TIME

READ 100, TIME

1((FORMAT(14.8)

NI=TOTAL NUMBER OF Y VALUES AT WHICH RESULTS ARE REQUIRED

READ 200,NH

YS=STARTING Y(SPACE VARIABLE); YF=FINAL Y; YC'S=VALUES OF Y'S AT WHICH RESULTS REQUIRED

READ 100, YS, YF, IYC(I), I=1, NH

IPR=10

F=1./NH

BETA=DLOG10(=DLOG(H))/DLOG10(YF)

PRINT 300, N, BETA, TIMEF, YF

300 FORMAT(1X, 'N = ', I3, 5X, 'BETA = ', F14.6, 5X, 'FINAL TIME = ', F14.6, 5X, 'Y-FINAL = ', E14.6, //)

R=NH

Z(1)=1.000

DO 400 I=2, N

Z(I)=Z(I-1)-H

Y(I)=(-DLOG(Z(I)))** (1./BETA)

DZ(I)=-BETA*(Y(I)**(BETA-1.))*Z(I)

400 DZ(I)=DZ(I)*((BETA-1.)/Y(I)-BETA*Y(I)**(BETA-1.))

TIME=0.000

CALL INTR(C, Y, NH)

DT=TIMEF/1000.+0.1

DI=IDT

500 TIME=TIME+DT

DO 600 I=2, N

DZZ=DZ(I)

DZZZ=DZZ(I)

YI=Y(I)

F1Y=F1(TIME, YI)

F2Y=F2(TIME, YI)

A(I-1)=2./DI+2.*DZZ*DZZ/F/R*F1Y-F3(TIME, YI)

T1=F1Y*DZZZ*DZZZ/F/R

T2=(F1Y*DZZZ+F2Y*DZZZ)/2./R

S(I-1)=-T1-T2

E(I-1)=-T1+T2

F1TY=F1(TIME-DT, YI)

F2TY=F2(TIME-DT, YI)

F3TY=F3(TIME-DT, YI)

F4TY=F4(TIME-DT, YI)

D(I-1)=F1TY*EZZ*DZZZ/R/F*(C(I+1)-2.*C(I)

+C(I-1))+ (C(I+1)-C(I-1))/2./F*(F1TY*(DZZZ+F2TY*DZZZ)

+F3TY*C(I)+F4TY+F4(TIME, YI))+2./DT*C(I)

600 CONTINUE

B(I)=D(I)-E(I)*G1(TIME)

D(N+1)=D(N)=B(N)*G2(TIME)

DO 700 I=1, NH

700 E(I)=E(I+1)

800 CALL TRID(NH, F, B, A, D)

C(1)=G1(TIME)

C(NH)=G2(TIME)

DO 900 I=2, N

```

C C C (I) = F (I-1)
  IP = IP + 1
  IF (IP.NE.IPR) GO TO 500
  YF = C
  LN = LN + 1
  CALL INTFLN(Y,YC,C,DENS,PETA)
  TIME S(LN) = TIME
  IF (TIME.GT.TIMEF) GO TO 1000
  GO TO 500
1000 CALL WRITE(YC,DENS,TIME S)
  STOP
  END

```

```

C SUBROUTINE INTFLN(Y,YC,C,DENS,PETA)

```

```

C *****
C THE SUBROUTINE INTERPOLATES DENS VALUES AT GIVEN YCS USING A
C THIRD DEGREE P(LYACNIA)
C *****

```

```

IMPLICIT REAL*8(A-H,I-Z)

```

```

REAL * 4 DENS

```

```

REAL * 4 HSS

```

```

DIMENSION Y(1000),YC(100),C(1000),DENS(100,200)

```

```

COMMON NM,NH,LM,NL,N

```

```

DO 6 J=1,NH

```

```

  DO 8 I=1,NL

```

```

    IF(YC(J).EQ.0.D00) GO TO 5

```

```

    IF(YC(J).GT.Y(I)) GO TO 8

```

```

    M=I

```

```

    GO TO 7

```

```

  5 CONTINUE

```

```

    M=Nl

```

```

  7 IF(N.LT.3) M=3

```

```

    IF(M.GT.NL-2) M=NL-1

```

```

    >=DEXP(-Y(N-2)*PETA)

```

```

    D1=C(M-1)-C(N-2)

```

```

    D2=C(M)-2.*C(N-1)+C(N-2)

```

```

    D3=C(N+1)-3.*C(M)+3.*C(N-1)-C(M-2)

```

```

    >>=DEXP(-YC(J)*PETA)

```

```

    (1->)->

```

```

    HSS=C(N-2)+UU*D1+UU*(UU-1.)/2.*D2+UU*(UU-1.)*(UU-2.)/6.*D3

```

```

    DENS(J,LM)=HSS

```

```

  GO TO 6

```

```

  8 DENS(J,LM)=C(I)

```

```

  6 CONTINUE

```

```

  RETURN

```

```

  END

```

```

SUBROUTINE YAIT (C,Y,NL)

```

```

*****
THE SUBROUTINE INITIALIZES THE GIVEN FUNCTION USING ITS VALUE
DEFINED FOR THE BOUNDARY CONDITION AT TIME=0
*****
IMPLICIT REAL*8 (A-H,C-Z)
DIMENSION C(1000), Y(1000)
DO 3 I= 1, N1
3 C(I)=G2(Y(I))
RETURN
END

```

```

SUBROUTINE TRID(A,C,B,A,C)
*****
THE SUBROUTINE SOLVES THE TRIDIAGONAL MATRIX BY CAUSSIAN
ELIMINATION (A METHOD)
*****
IMPLICIT REAL*8 (A-H,C-Z)
DIMENSION C(1), B(1), A(1), D(1)
NN=N-1
DO 1 I=1, NN
D(I)=D(I)/A(I)
B(I)=B(I)/A(I)
T1=D(I)
T2=B(I)
IF(DABS(T1).LT.1.E-30)T1=0.00
IF(DABS(T2).LT.1.E-30)T2=0.00
C(I+1)=C(I+1)-T1*B(I)
1C A(I+1)=A(I+1)-T2*A(I)
IF(DABS(D(N)).LT.1.E-30)D(N)=0.00
C(N)=C(N)/A(N)
D(2*1)=1, NN
C(N-1)=D(N-1)-B(N-1)*C(N-1+1)
IF(DABS(C(N-1)).LT.1.E-30)C(N-1)=0.00
2C CONTINUE
RETURN
END

```

```

SUBROUTINE WRITE(YC,CEAS,TIMES)
*****
MULTI SUBROUTINE : IT CAN PLOT AND/OR PRINT AS SPECIFIED
BY THE USER THROUGH THE VALUE OF K2 IN THE MAIN PROGRAM
*****
IMPLICIT REAL*8 (A-H,C-Z)
REAL*4 YC
DIMENSION YC(1000), CEAS(100,200), TIMES(1000)
COMMON NN, NH, LM, NI, N, K1, K2, K3, K4
YMAX=1.E-5C
YMIN=1.E-50
>NIN=TIMES(1)

```

```

XMAX=TIMES(I,N)
DO 2 J=1,NH
DO 1 I=1,1N
IF(DENS(J,I).GT.YMAX) YMAX=DENS(J,I)
1 IF(DENS(J,I).LT.YMIN) YMIN=DENS(J,I)
2 CONTINUE
IF(K2.EC.1) XXMAX=1./XMIN
IF(K2.EC.1) XXMIN=1./XMAX
IF(K2.EC.1) XMAX=XXMAX
IF(K2.EC.1) XMIN=XXMIN
IF(K1.EC.1) XMIN=DLOG10(XMIN)
IF(K1.EC.1) XMAX=DLOG10(XMAX)
IF(K1.EC.1) YMAX=DLOG10(YMAX)
IF(K1.EC.1.AND.YMIN.EC.0.C00) GO TO 20
IF(K1.EC.1.AND.YMIN.NE.0.C00) YMIN=DLOG10(YMIN)
10 HY=XMAX-XMIN
HY=YMAX-YMIN
WRITE(6,3) XMIN,XMAX,YMIN,YMAX
3 FORMAT(1X,'XMIN = ',E12.6,10X,'XMAX = ',E12.6,/,1X,'YMIN = ',
E12.6,10X,'YMAX = ',E12.6,/,1X,'YC,5 = ')
IF(K3.NE.1) GO TO 40
CALL FLCT(7)
CALL FLCT(12,XMIN,XMAX,10.,HX,YMIN,YMAX,10.,HY)
DO 5 J=1,NH
DO 4 I=1,1N
XX=TIMES(I)
YY=DENS(J,I)
IF(K2.EC.1) XX=1./TIMES(I)
IF(K1.EC.1) XX=DLOG10(XX)
IF(K1.EC.1) YY=DLOG10(YY)
CALL FLCT(90,XX,YY)
4 CONTINUE
CALL FLCT(99)
WRITE(6,6) YC(J)
6 FORMAT(10X,E12.6)
5 CONTINUE
IF(K4.NE.1) GO TO 11
40 WRITE(6,25) (YC(I),I=1,NH)
25 FORMAT(1X,'TIME PARAMETER',T20,'VALUE OF FUNCTION AT INTERPOLATED
CYC,5',/,1X,T19,' @ (1X,E12.6)')
DO 30 I=1,1N
IF(K2.EC.1) TIMES(I)=1./TIMES(I)
IF(K1.EC.1) TIMES(I)=DLOG10(TIMES(I))
DO 45 J=1,NH
45 IF(K1.EC.1) DENS(J,I)=DLOG10(DENS(J,I))
WRITE(6,35) TIMES(I),(DENS(J,I),J=1,NH)
35 FORMAT(1X,E12.6,T19,6(1X,E12.6))
30 CONTINUE
GO TO 11
20 WRITE(6,15)
15 FORMAT(1X,'DLOG10(YMIN) CANNOT BE TAKEN SINCE ARGUMENT = 0.0!')

```


11 RETURN

END

```
FUNCTION F1(TIME,YI)
IMPLICIT REAL*8(A-H,C-Z)
F1=1.0-C2
RETURN
END
```

```
FUNCTION F2(TIME,YI)
IMPLICIT REAL*8(A-H,C-Z)
F2=-1.0-D02
RETURN
END
```

```
FUNCTION F3(TIME,YI)
IMPLICIT REAL*8(A-H,C-Z)
F3=C.DC0
RETURN
END
```

```
FUNCTION F4(TIME,YI)
IMPLICIT REAL*8(A-H,C-Z)
F4=C.DC0
RETURN
END
```

```
FUNCTION G1(TIME)
IMPLICIT REAL*8(A-H,C-Z)
G1 = DSIN( 2* 3.14159265 / 1000. * TIME )
RETURN
END
```

```
FUNCTION G2(TIME)
IMPLICIT REAL*8(A-H,C-Z)
G2=C.DC0
RETURN
END
```

```
FUNCTION G3(TIME)
IMPLICIT REAL*8(A-H,C-Z)
G3=C.DC0
RETURN
END
```

APPENDIX 2

The solution of

$$\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2} - u \frac{\partial C}{\partial x} \quad (1)$$

subject to the initial and boundary conditions

$$C(x, 0) = 0; \quad C(0, t) = C_0 \sin \omega t; \quad C(\infty, t) = 0 \quad (2)$$

is given by

$$C(x, t) = \frac{C_0}{2} \exp \left\{ \frac{ux}{2D} - \mu t \right\} \left[\exp(\varphi_1) \sin \varphi_2 \right. \\ \left. \{ 1 - \operatorname{Re} \{ \operatorname{erf}(z_+) \} \} + \exp(\varphi_3) \sin \varphi_4 \{ 1 - \operatorname{Re} \{ \operatorname{erf}(z_-) \} \} \right] \\ - \left\{ \exp(\varphi_1) \cos \varphi_2 \operatorname{Im} \{ \operatorname{erf}(z_+) \} + \exp(\varphi_3) \cos \varphi_4 \right. \\ \left. \operatorname{Im} \{ \operatorname{erf}(z_-) \} \right] \quad (3)$$

where

$$\begin{aligned} \varphi_1 &= \operatorname{tr} \cos \theta + x (r/D)^{1/2} \cos \theta/2 \quad \} \\ \varphi_2 &= \operatorname{tr} \sin \theta + x (r/D)^{1/2} \sin \theta/2 \quad \} \\ \varphi_3 &= \operatorname{tr} \cos \theta - x (r/D)^{1/2} \cos \theta/2 \quad \} \\ \varphi_4 &= \operatorname{tr} \sin \theta - x (r/D)^{1/2} \sin \theta/2 \quad \} \end{aligned} \quad (4)$$

$$r = (\mu^2 + \omega^2)^{1/2}; \quad \tan \theta = \omega/\mu; \quad \mu = u^2/4D$$

Infinite series approximation for complex error function is

$$\operatorname{erf}(x + iy) = \operatorname{erf}(x) + \frac{e^{-x^2}}{2\pi x} \left[(1 - \cos 2xy) \right.$$

$$\left. + i \sin 2xy \right] + \frac{2}{\pi} e^{-x^2} \sum_{n=1}^{\infty} \frac{e^{-n^2/4}}{n^2 + 4x^2} \left[f_n(x, y) \right.$$

$$\left. + i g_n(x, y) \right] + \epsilon(x, y)$$

where

$$f_n(x, y) = 2x - 2x \cosh ny \cos 2xy + n \sinh ny \sin 2xy$$

$$g_n(x, y) = 2x \cosh ny \sin 2xy + n \sinh ny \cos 2xy$$

$$\text{and } \epsilon(x, y) \approx 10^{-16} \left| \operatorname{erf}(x + iy) \right|$$

where z_{\pm} are defined as

$$z_{\pm} = \left. \begin{aligned} & x/2 (Dt)^{1/2} \pm (rt)^{1/2} \cos \theta/2 \pm iy \\ & \end{aligned} \right\} \quad (5)$$

$$\text{and } y = (rt)^{1/2} \sin \theta/2$$

Various terms appearing in equations (3) - (5) may be expressed in terms of a dimensionless distance parameter, (xu/D) , and the dispersion time constant, (D/u^2) . For example,

$$(i) \quad ux/2D - \mu t \text{ appearing in eqn. (3) may be written as } ux/2D - t/(4D/u^2)$$

$$(ii) \quad \phi_1 = tr \cos \theta + x (r/D)^{1/2} \cos \theta/2 \\ = t\mu(1 + \omega^2/\mu^2)^{1/2} + (xu/2D) (1 + \omega^2/\mu^2)^{1/4}$$

and similar expressions for φ_2 , φ_3 and φ_4 .

$$(iii) z_+ = (xu/2D) (tD/u^2)^{-1/2} + (t/(2D/u^2))^{1/2}$$

$$\{1 + \omega^2/(4D/u^2)^2\} \cos \theta/2$$

$$+ i\{(t/(4D/u^2))^{1/2} \{1 + \omega^2/(4D/u^2)^2\}^{1/4} \sin \theta/2.$$