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Short Communication

Infrared attenuation due to phase change from amorphous to crystalline observed in astrochemical propargyl ether ices

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ABSTRACT

Astrochemical ices are known to undergo morphological changes, from amorphous to crystalline, upon warming the ice from lower (10 K) to higher temperatures. Phase changes are mostly identified by the observation of significant changes in the InfraRed (IR) spectrum, where the IR bands that are broad in the amorphous phase are narrower and split when the ice turns crystalline. To-date all the molecules that are studied under astrochemical conditions are observed to follow such a behaviour without significant attenuation in the IR wavelength. However, in this paper we report a new observation when propargyl ether ($C_3H_3OC_3H_3$) is warmed from the amorphous phase, at 10 K, through the phase transition temperature of 170 K, the crystalline ice being found to strongly attenuate IR photons at the mid-IR wavelengths.

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1. Introduction

InfraRed (IR) spectroscopy plays a vital role in the study of interstellar molecules both in space-based observations and in laboratory simulations. The *Infrared Space Observatory* and *Spitzer* space telescope operating in the Mid and Far IR wavelength has been widely used to explore the icy nature of molecular clouds and their chemical composition by comparing the spectral features observed in the InterStellar Medium (ISM) with measurements in laboratory ice analogues [1–4].

While investigating astrochemical ices and their corresponding IR spectra at temperatures compatible to the ISM (10 K) the laboratory ices are prepared in what is usually attributed to the amorphous phase. After deposition, such ices may be warmed gradually or at a fixed heating rate (such as 0.5 to $few\ K\ min^{-1}$) to mimic heating of such ices in the ISM e.g. by shocks. Subtle changes appearing in the IR spectra during the warmup phase may be related to the conformational changes related to homodimers [5]. However, significant changes that appear in the spectra, such as band splitting along with the appearance of sharper bands, are related to the phase change from amorphous to

crystalline that take place within the ice. Recently, by observing such changes in the IR spectra an unexpected reversible phase change was reported to occur in astrochemical ethanethiol ices [6].

In the present study we report a temperature dependent study of the IR spectrum of propargyl ether ($C_3H_3OC_3H_3$), an isomer of phenol (C_6H_5OH) which is reported (tentative detection with a predicted column density of $\sim 8 \pm 4 \times 10^{14}$) to be present in the Orion KL [7] using the IRAM 30 m telescope. Indeed, propargyl ether is regarded as one of the interesting molecules, awaiting discovery in ISM, in understanding the chemical complexity that surrounds the evolution of aromatic molecules in icy mantles. Previous laboratory experiments on propargyl containing compounds are limited to the study of propargyl alcohol ices [8–10] and to-date no IR and Vacuum UltraViolet (VUV) spectroscopic data are available for propargyl ether ices at temperatures commensurate with the ISM. Accordingly, propargyl ether IR and VUV spectra were recorded as a function of temperature and the observations made in both the amorphous and crystalline phases of propargyl ether.

2. Experimental methodology

Experiments were carried out in the 10 K astrochemical ices simulation chamber housed at the Physical Research Laboratory (PRL), India. The apparatus consists of an UltraHigh Vacuum (UHV) chamber operating at a base pressure of about 10^{-10} mbar. A sample holder containing

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a Zinc Selenide (ZnSe) window, used as a dust analog to form icy mantles, was cooled to 10 K using a closed cycle cryostat. Propargyl ether molecules drawn from a liquid reservoir were allowed to deposit, 30 mbar in 230 s, on the cooled ZnSe window to form an ice film. For deposition at 150 K and 170 K, a similar deposition rate was maintained. A Fourier Transform InfraRed (FTIR) spectrometer operating in the mid-IR (4000–550 cm^{-1} , step size 2 cm^{-1}) is used to probe the morphology of the propargyl ether ices. Spectra of the ice were then recorded as the ice was warmed to higher temperatures, at a rate of 5 K min^{-1} , until the propargyl ether sublimated from the ZnSe window.

Vacuum UltraViolet (VUV) photoabsorption spectra of propargyl ether ices were also recorded using the experimental setup at the O3A1 beamline facility at NSRRC [11]. A Lithium Fluoride (LiF) window was used as substrate, cooled down to 10 K onto which propargyl ether was deposited to form an ice layer. In order to obtain the VUV spectra of such films a spectrum of the LiF window was obtained prior to deposition *Io* and after deposition, *I*. By employing the Beer-Lambert law the absorbance of the ice could then be obtained. By warming the sample to higher temperatures and recording spectra at each temperature the temperature dependent VUV spectra of propargyl ether ices were obtained. In both the experiments the propargyl ether samples (purity 98%) were obtained from Sigma Aldrich.

3. Results and discussion

The IR spectrum recorded after depositing propargyl ether molecules at 10 K, was observed to contain several peaks in the 4000–550 cm^{-1} mid-IR region (Table 1; Fig. 1). A peak corresponding to CH stretching was observed at 3492 cm^{-1} and the band at 3061 cm^{-1} was assigned to the characteristic CH_2 asymmetric stretching vibration. Whereas, the band at 2964 cm^{-1} was assigned to the CH_2 symmetric stretching. The $\text{C}\equiv\text{C}$ stretching vibration in the propargyl ether molecule was observed at 2120 cm^{-1} . The set of bands observed at 1445 cm^{-1} , 1404 cm^{-1} /1347.5 cm^{-1} and 1267 cm^{-1} /1248 cm^{-1} was attributed to the CH_2 scissoring, CH_2 wagging and CH_2 twisting vibrations, respectively. The band at 1080.5 cm^{-1} was then attributed to the C—O stretching, whilst the adjacent band at 1035.6 cm^{-1} corresponds to the CH_2 rocking vibration. The ether link, in the propargyl ether molecule, COC bending vibration was observed at 1007 cm^{-1} . However, this band can also be attributed to C—C stretching vibration too. The bands at 980 cm^{-1} and 942 cm^{-1} can be attributed to the CH_2 rocking/C—O stretching and C—C stretching vibrations, respectively. The band closest to the spectral limit used in this experiment was observed at 610 cm^{-1} and that corresponds to the CH bending vibration. The band at 578.5 cm^{-1} can be attributed to the CH torsion. Combination bands are also provided in Table 1. Moreover, the full list of the calculated frequencies are provided as supplementary material.

After slowly warming the ice, the spectrum recorded at 150 K was observed to have become more intense when compared to that at 10 K though the spectral features are similar. In addition to that there was significant reduction in the transmittance (Fig. 1). This is peculiar since the number of absorbers has not changed. At still higher temperature the spectrum recorded at 170 K changes dramatically with significant alteration in the spectra especially in the 1500–600 cm^{-1} region (Fig. 1) and the transmittance was observed to have reduced significantly. It was not straight forward to assign the bands, in the 1500–600 cm^{-1} , in the IR spectrum recorded at 170 K, so the bands observed at 170 K are given in a separate table (Table 2). Apart from the sharp spectral changes the attenuation observed in the IR wavelength is seen to increase at the higher temperature as the phase changes from amorphous to crystalline in the propargyl ether ices. This is an extremely unexpected result. Such a dramatic change has consequences for interpreting astronomical data where such a change in absorption could be interpreted as a change in number density.

Table 1

Peak position of IR bands observed in amorphous propargyl ether ice formed at 10 K compared with calculated band positions and corresponding assignments.

Amorphous (10 K), cm^{-1}	Calculated, cm^{-1}	Assignments (mode)
3960	3972.2 3969.4	($\nu_{20} + \nu_5$) ($\nu_{20} + \nu_6$)
3914.5		
3627	3626.9 3645.1	($\nu_{25} + \nu_3$) ($\nu_{22} + \nu_6$)
3492	3467.2 3467.3	CH stretching--(ν_2/ν_1)
3283	3287.1 3284.3	($\nu_{28} + \nu_5$) ($\nu_{28} + \nu_6$)
3248.5	3266.5 3263.7	($\nu_{29} + \nu_5$) ($\nu_{29} + \nu_6$)
3061	3081.6 3086.8	CH_2 asymmetric stretching --(ν_4) CH_2 asymmetric stretching --(ν_3)
2964	3034.1 3031.3	CH_2 symmetric stretching --(ν_5) CH_2 symmetric stretching --(ν_6)
2908		
2860	2837.7	($\nu_{12} + \nu_9$)
2120	2121.0 2119.9	$\text{C}\equiv\text{C}$ stretching --(ν_7) $\text{C}\equiv\text{C}$ stretching --(ν_8)
2021	2036.1	($\nu_{21} + \nu_{11}$)
1953.2	1973.9	($\nu_{21} + \nu_{12}$)
1624	1620.5	($\nu_{21} + \nu_{18}$)
1445	1487.4 1476.5	CH_2 scissoring--(ν_9) CH_2 scissoring--(ν_{10})
1404	1412.5	CH_2 wagging --(ν_{11})
1383	1391.4	($\nu_{30} + \nu_{14}$)
1372	1365.1	($\nu_{26} + \nu_{20}$)
1361.7	1347.9	($\nu_{28} + \nu_{15}$)
1347.5	1350.3	CH_2 wagging--(ν_{12})
1293.5	1290.3	($\nu_{28} + \nu_{16}$)
1267	1251.8	CH_2 twisting --(ν_{13})
1248	1243.3	CH_2 twisting --(ν_{14})
1143.5	1145.0	($\nu_{30} + \nu_{18}$)
1080.5	1094.9	C—O stretching --(ν_{15})
1047	1040.8	($\nu_{26} + \nu_{22}$)
1035.6	1037.3	CH_2 rocking (ν_{16})
1007	1010.5	C—C stretching/COC bending/ CH_2 rocking (ν_{17})
980	994.5 996.9	CH_2 rocking --(ν_{19}) C—O stretching --(ν_{18})
942	938.0	C—C stretching --(ν_{20})
918	931.6	($\nu_{27} + \nu_{21}$)
903.5	921.8	($\nu_{27} + \nu_{22}$)
881.5	884.9 883.0	($\nu_{27} + \nu_{23}$) ($\nu_{27} + \nu_{24}$)
744	761.8 724.9 723.0	($\nu_{30} + \nu_{22}$) ($\nu_{30} + \nu_{23}$) ($\nu_{30} + \nu_{24}$)
685.5	680.0	($\nu_{28} + \nu_{26}$)
646.5	657.4	($\nu_{32} + \nu_{23}$)
610	623.5 613.7	CH bending --(ν_{21}) CH bending --(ν_{22})
range out of 4000–600	576.8 574.9 575.1	CH torsion --(ν_{23}) CH torsion --(ν_{24}) ($\nu_{30} + \nu_{26}$)
	540.1	CCO in plane bending --(ν_{25})
	427.0	CCC in plane bending --(ν_{26})
	308.1	CCC in plane bending --(ν_{27})
	252.9	CCC out of plane bending --(ν_{28})
	232.3	CCC out of plane bending --(ν_{29})
	148.0	CCH bending --(ν_{30})
	95.1	COC out of plane bending --(ν_{31})
	80.6	CCH bending --(ν_{32})
	43.7	CH_2 bending --(ν_{33})

Note: Vibrational frequencies of propargyl ether were computed using Gaussian 09 program [12]. We ran frequency calculation after optimization at MP2/aug-cc-pVDZ level of theory. Vibrational wavenumbers of harmonic fundamental and combination bands of propargyl ether are presented above. We are providing the calculated datafile as a Supplementary material of this article.

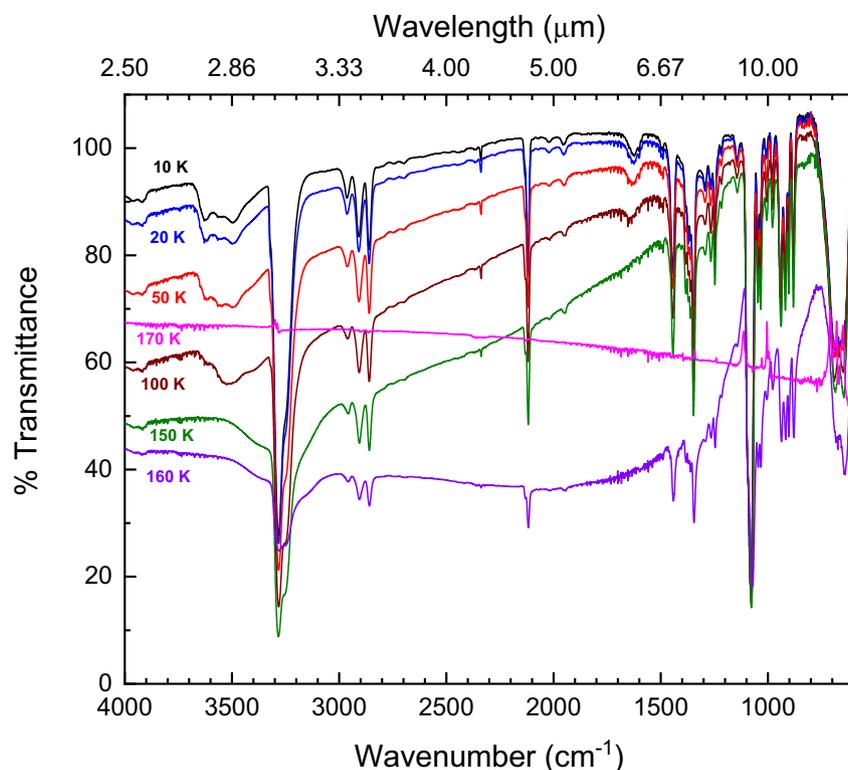


Fig. 1. IR spectra of propargyl ether ice as deposited at 10 K and warmed at the rate of 5 K min⁻¹ to higher temperatures, 170 K (phase change).

Table 2

Tentative peak positions of bands observed in crystalline propargyl ether ice obtained by warming the ice from 10 K to 170 K.

Crystalline (170 K), cm ⁻¹		
4338	2054.5	1269.3
3945.5	2009.6	1236.7
3913.7	1977.5	1094
3291.7	1457.6	1026.1
3283	1451.8	1000
2924	1444	987
2900.8	1403.7	937
2864.7	1395.5	913.4
2785	1350.2	701.4
2130	1336.7	686.5
2125.4	1321	673
2106	1296.3	635
2071.8		

The spectrum recorded after cooling the crystalline propargyl ether ice to 10 K, was observed to have similar transmittance as to the crystalline ice at 170 K (Fig. 2). The phase change was irreversible and the transmittance observed was found to be related to the phase change from amorphous to crystalline. Similar drastic reduction in transmittance was observed in the IR spectrum recorded after depositing propargyl ether samples at 170 K (Fig. 2).

Two prominent bands were observed in the VUV spectra of propargyl ether recorded at 10 K (i) in the 170–210 nm with a peak at 192 nm and (ii) in the 110–170 nm region with a broad band 130–160 nm re-

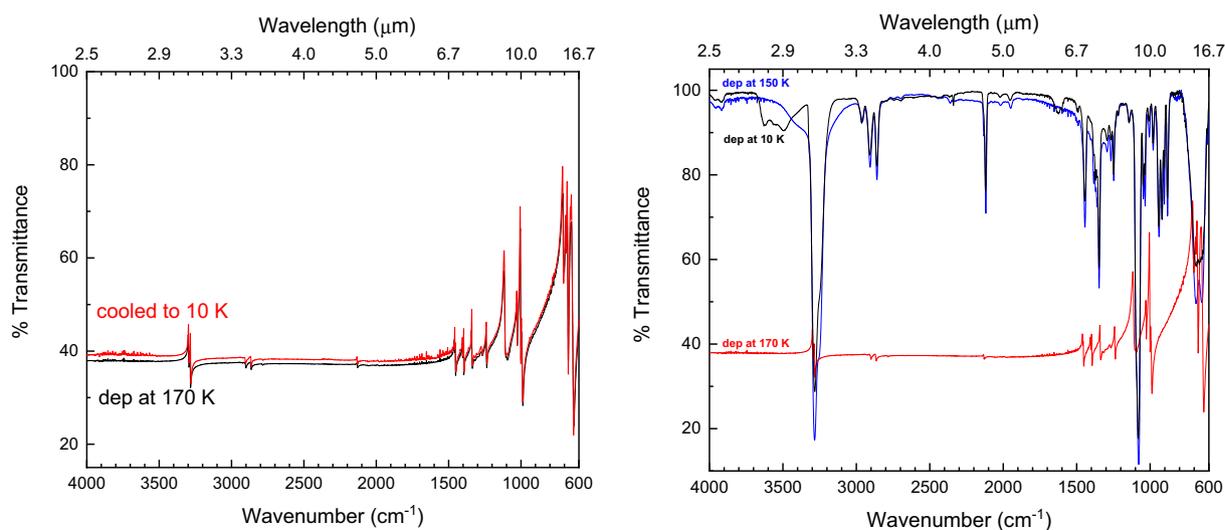


Fig. 2. (Left) IR spectra of crystalline propargyl ether ice deposited at 170 K compared with the crystalline ice spectrum recorded at 10 K by cooling the sample and (right) IR spectra of propargyl ether ices deposited at 10 K, 150 K and 170 K.

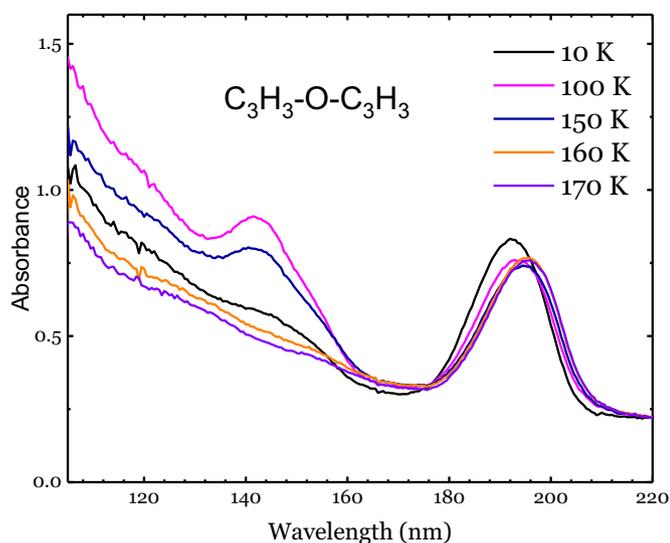


Fig. 3. The VUV spectra of propargyl ether ices as a function of temperature.

gion (Fig. 3). Upon warming the ice to higher temperatures, 100 K–150 K the band at 130–160 nm was observed to grow in intensity and the peak at 192 nm seen to red shift by ~2 nm (Fig. 3). The spectrum at 170 K was observed to have spectral features in both the 110–170 nm and 170–210 nm regions. Further VUV spectral data analysis is not detailed in this manuscript, because the motivation behind the VUV spectral recording of propargyl alcohol was to investigate whether attenuation happens at wavelengths other than IR when propargyl ether ice turns from amorphous to crystalline. The VUV spectra recorded as a function of temperature clearly indicate no significant attenuation of propargyl ether ice in the 110–210 nm wavelength.

4. Conclusion

Propargyl ether ices formed under astrochemical icy conditions were probed by IR and VUV spectroscopic techniques. IR spectroscopy revealed that after the phase change occurred, from amorphous to crystalline, IR photons are strongly attenuated in the crystalline propargyl ether ices. For the ice probed in the VUV spectral region such an attenuation linked to phase change was not observed. Therefore, here we

report the first astrochemical ice, propargyl ether, that attenuates IR photons in the mid-IR region upon phase change from amorphous to crystalline. Such an effect may be due to change in refractive index of the propargyl ether ice.

This experimental result has implications as the next space borne and highly sensitive, James Webb space telescope, which will be probing the ISM at IR wavelengths. Due to the presence of such astrochemical ices attenuating IR photons, interpretation and analysis of IR data may not be straight forward. Therefore, this requires more laboratory experiments to find if such effects are observed in other pure astrochemical ices or in ice mixtures.

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.saa.2019.117393>.

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